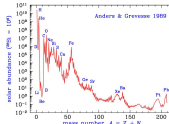
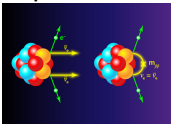
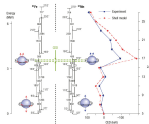


# NUCLEAR STRUCTURE WITHIN THE SHELL MODEL: NEW FRONTIERS

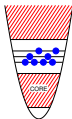
Dao Duy Duc and Frédéric Nowacki  
(IPHC-Strasbourg)

GANIL - Autrans Méaudre en Vercors  
September 29th, 2021



# 1 Shell model framework

## Shell model description of the nucleus



$$\mathcal{H}_{\text{eff}}\Psi_{\text{eff}} = E\Psi_{\text{eff}} \quad (\text{N. Smirnova's talk this morning})$$

$$\mathcal{H}_{\text{eff}} = \mathcal{H}_m + \mathcal{H}_M$$

### Exact Diagonalization

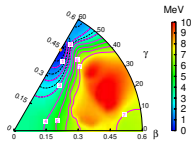
- Valence spaces
- Effective interactions
- Gigantic dimension:  $D \sim \binom{d_\pi}{p} \cdot \binom{d_\nu}{n}$

(Strasbourg codes **ANTOINE**, **NATHAN**)

### Generator Coordinate Method (GCM)

- ▷ Deformed Hartree-Fock:

$$|\Phi\rangle = \prod_{i=1}^A a_i^\dagger |0\rangle$$



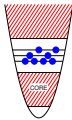
- ▷ Rotational symmetry restoration:  $\mathcal{P}_{MK}^J |\Phi(q)\rangle \propto |J, M\rangle$
- ▷ Mixing of non-orthogonal states by the GCM:

$$|\Psi^{(J,M)}\rangle = \sum f^{(J)}(q, K) \mathcal{P}_{MK}^J |\Phi(q)\rangle$$

$$\sum_{q,K} \left( H_{K'K}^{(J)}(q', q) - E^{(J)} N_{K'K}^{(J)}(q', q) \right) f^{(J)}(q, K) = 0$$

# 1 Shell model framework

## Shell model description of the nucleus



$$\mathcal{H}_{\text{eff}}\Psi_{\text{eff}} = E\Psi_{\text{eff}} \quad (\text{N. Smirnova's talk this morning})$$
$$\mathcal{H}_{\text{eff}} = \mathcal{H}_m + \mathcal{H}_M$$

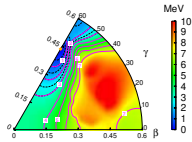
### Exact Diagonalization

- Valence spaces
- Effective interactions
- Gigantic dimension:  $D \sim \binom{d_\pi}{p} \cdot \binom{d_\nu}{n}$

### Generator Coordinate Method (GCM)

- ▷ Deformed Hartree-Fock:

$$|\Phi\rangle = \prod_{i=1}^A a_i^\dagger |0\rangle$$

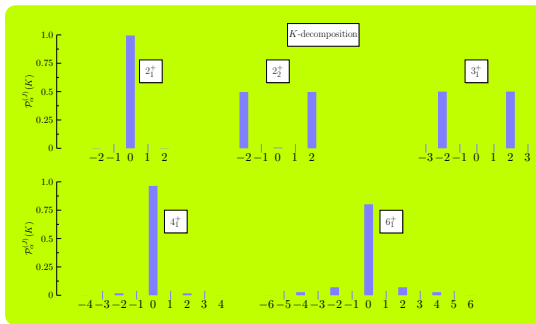
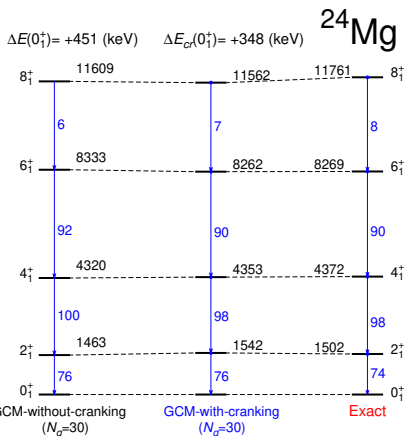


- ▷ Rotational symmetry restoration:  $\mathcal{P}_{MK}^J |\Phi(q)\rangle \propto |J, M\rangle$

### Problem of selecting relevant deformed states: **Minimization Technique**

*"The first point is the one such that it minimizes the energy. The second point is chosen in such a way that the energy obtained from diagonalizing the Hamiltonian in the 2-dimensional space be a minimum. One proceeds in the same way to determine the third basis vector etc..."* (E. Caurier, Proc. on GCM, BLG report **484** (1975), Bouten, M., Van Leuven, P. (ed.))

# 2 Systematics in sd nuclei with the USDB interaction



$^{24}\text{Mg}$

$E(2_2^+)$  (keV)

$E(3_1^+)$

$B(E2; 2_2^+ \rightarrow 2_1^+)$  ( $e^2\text{fm}^4$ )

$B(E2; 3_1^+ \rightarrow 2_2^+)$

Exact

4116

5070

14

131

$(\beta, \gamma)$

4059

5140

16

138

$(\beta, \gamma, \langle \hat{J}_z \rangle)$

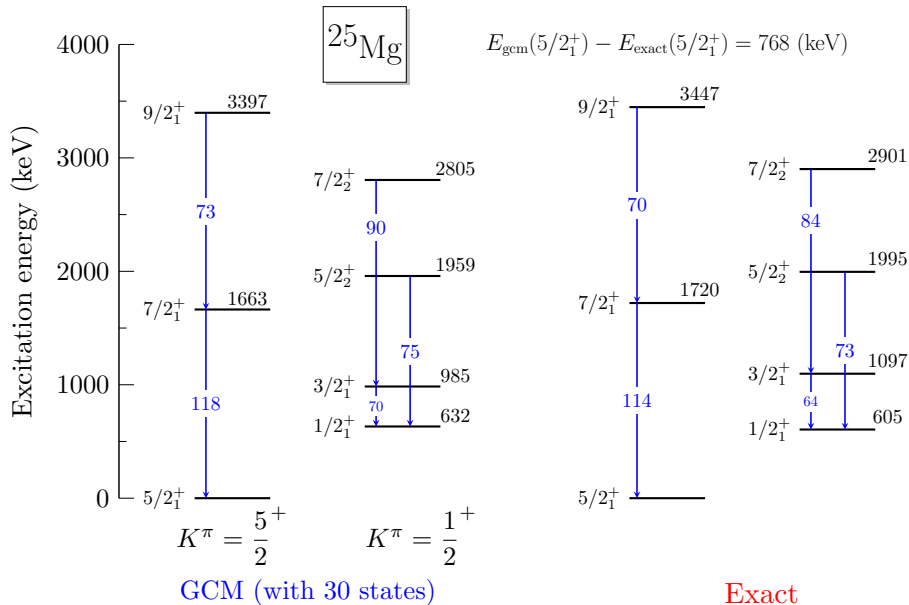
4121

5199

17

137

# 4 Systematics in sd nuclei with the USDB interaction



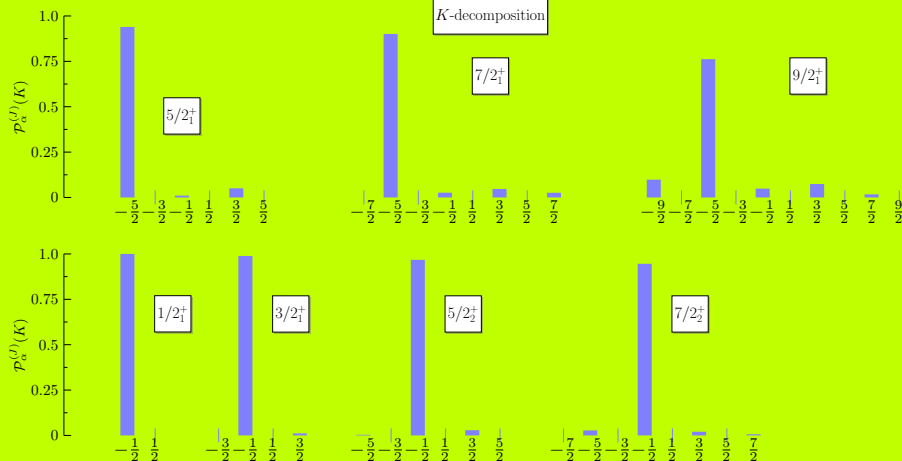
# 4 Systematics in sd nuclei with the USDB interaction

4000

$^{25}\text{Mg}$

$$E_{\text{gcm}}(5/2_1^+) - E_{\text{exact}}(5/2_1^+) = 768 \text{ (keV)}$$

3117



# 4 Systematics in sd nuclei with the USDB interaction

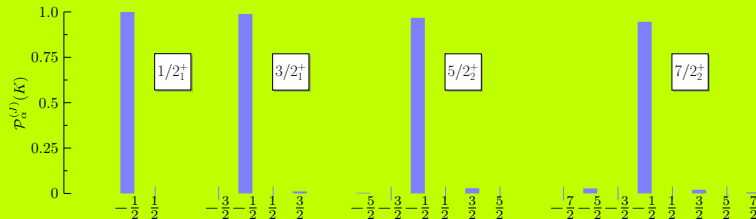
4000  $\uparrow$



$$E_{\text{gcm}}(5/2_1^+) - E_{\text{exact}}(5/2_1^+) = 768 \text{ (keV)}$$

## Summary for this part:

- Efficient method to choose basis states in the GCM
- Nature of the GCM state: K-decomposition complementary to shell model interpretation in spherical orbits
- Monopole/Multipole decomposition of the effective Hamiltonian

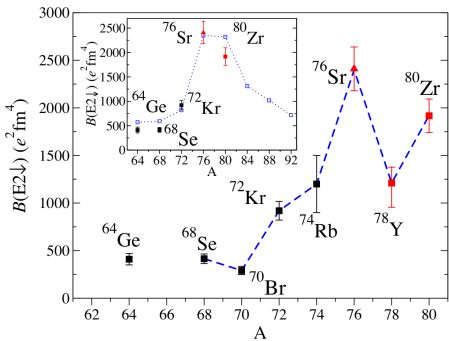


# 5 Collectivity in heavy $N = Z$ nuclei

PHYSICAL REVIEW LETTERS **124**, 152501 (2020)

## Establishing the Maximum Collectivity in Highly Deformed $N=Z$ Nuclei

R. D. O. Llewellyn,<sup>1,\*</sup> M. A. Bentley,<sup>1</sup> R. Wadsworth,<sup>1,†</sup> H. Iwasaki,<sup>2,3</sup> J. Dobaczewski,<sup>1,4</sup> G. de Angelis,<sup>5</sup> J. Ash,<sup>2,3</sup>  
 D. Bazin,<sup>2,3</sup> P. C. Bender,<sup>2,‡</sup> B. Cederwall,<sup>6</sup> B. P. Crider,<sup>2,§</sup> M. Doncel,<sup>7</sup> R. Elder,<sup>2,3</sup> B. Elman,<sup>2,3</sup> A. Gade,<sup>2,3</sup>  
 M. Grinder,<sup>2,3</sup> T. Haylett,<sup>1</sup> D. G. Jenkins,<sup>1</sup> I. Y. Lee,<sup>8</sup> B. Longfellow,<sup>2,3</sup> E. Lunderberg,<sup>2,3</sup> T. Mijatović,<sup>2,||</sup>  
 S. A. Milne,<sup>1</sup> D. Muir,<sup>1</sup> A. Pastore,<sup>1</sup> D. Rhodes,<sup>2,3</sup> and D. Weisshaar<sup>2</sup>



- Presence of  $d5/2 \Rightarrow$  deformation
- LNPS matrix elements

Same space for  $\pi, \nu$

g9/2

p1/2

f5/2

p3/2

JUN45

d5/2

g9/2

p1/2

f5/2

p3/2

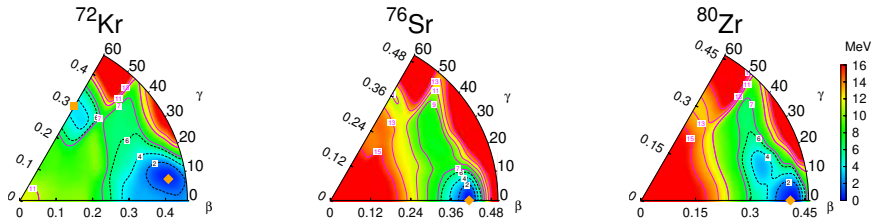
JUN45.LNPS

<sup>56</sup><sub>28</sub>Ni<sub>28</sub>

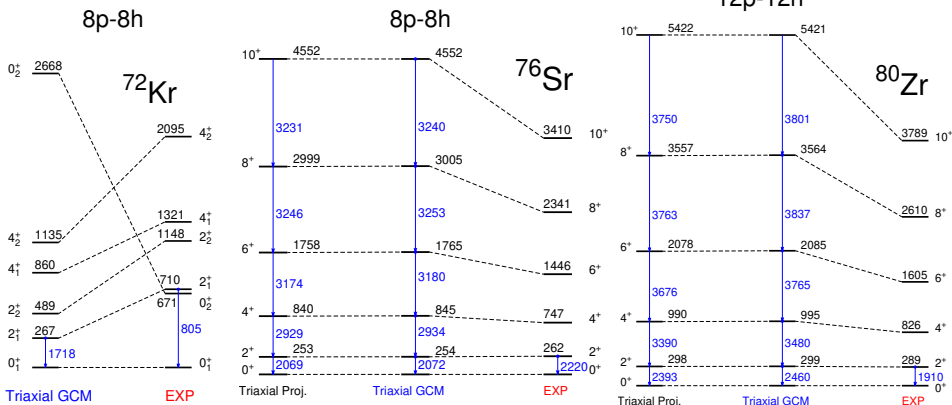
FIG. 3. Schematics of the  $B(E2\downarrow)$  values for the  $N = Z$  nuclei



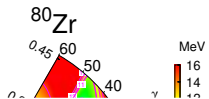
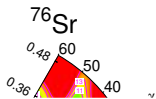
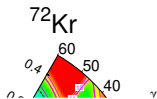
# 6 Collectivity in heavy $N = Z$ nuclei



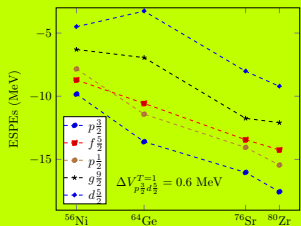
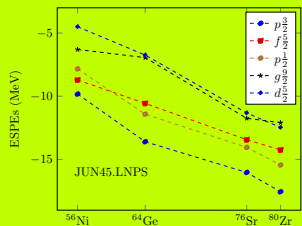
- HF minimum and ground state structures:



# 6 Collectivity in heavy $N = Z$ nuclei

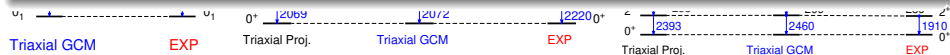


## Effective Single Particle Energies (ESPEs)



## Summary for this part:

- ◆ First calculations of heavy  $N = Z$  nuclei within the Shell Model
- ◆ Good description in the case of  $^{76}\text{Sr}$ ,  $^{80}\text{Zr}$
- ◆ Work under progress to determine the location of the  $d_{5/2}^+$  orbit (systematics of second excited states  $0_2^+$  where JUN45 becomes inadapted)

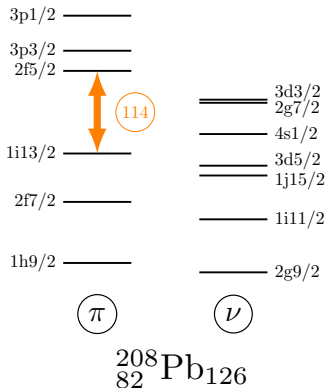


### Kuo-Herling interaction:

- <sup>208</sup>/<sub>82</sub>Pb<sub>126</sub> core, realistic TBMEs
- $82 \leq Z \leq 126$  shells for proton and  $126 \leq N \leq 184$  for neutrons
- monopole corrections (3N force)

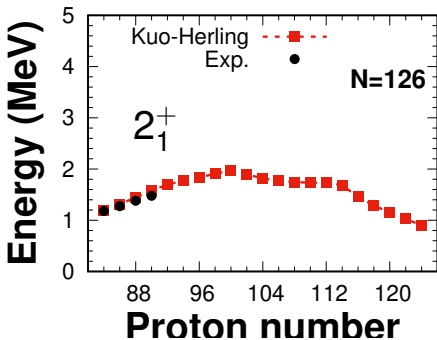
*E. Caurier and F. Nowacki,*

*PRL 87 (2001),072511*



### Calculations: NATHAN & DFSM

- ◊ Diagonalization within the seniority scheme along the chains of  $N = 126$  and  $N = 184$
- ◊ Deformed Hartree-Fock and Angular momentum projection plus shapes mixing through the GCM: <sup>254</sup>No

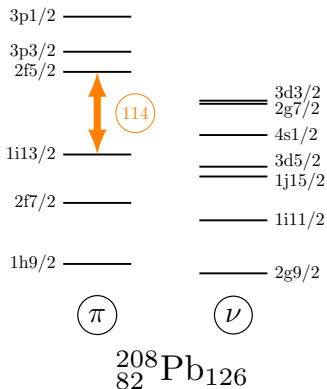


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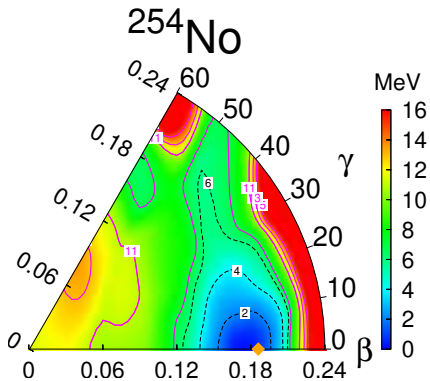
*E. Caurier and F. Nowacki,*

*PRL 87 (2001),072511*

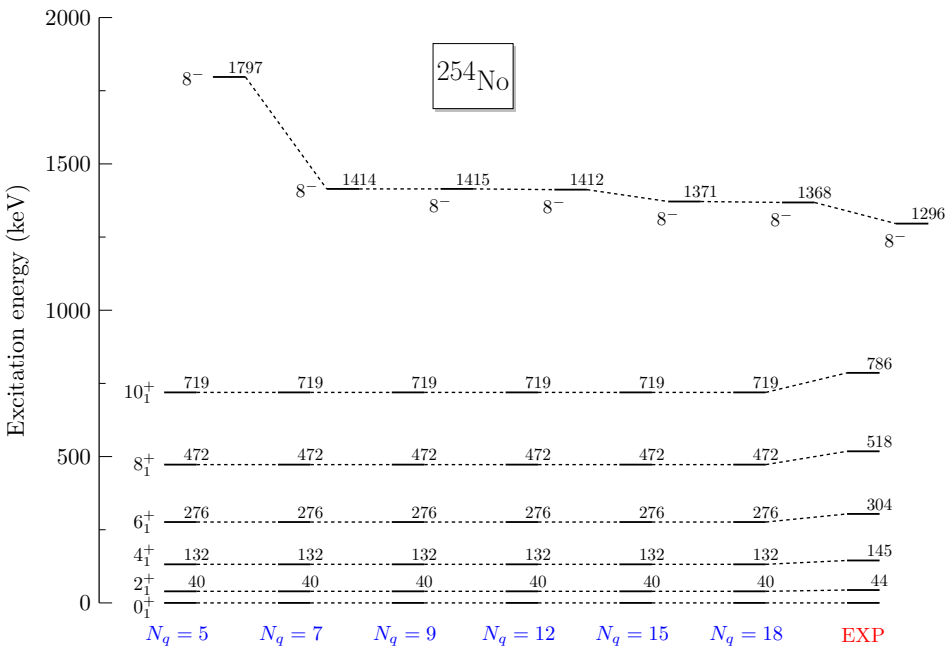


### Calculations: NATHAN & DFSM

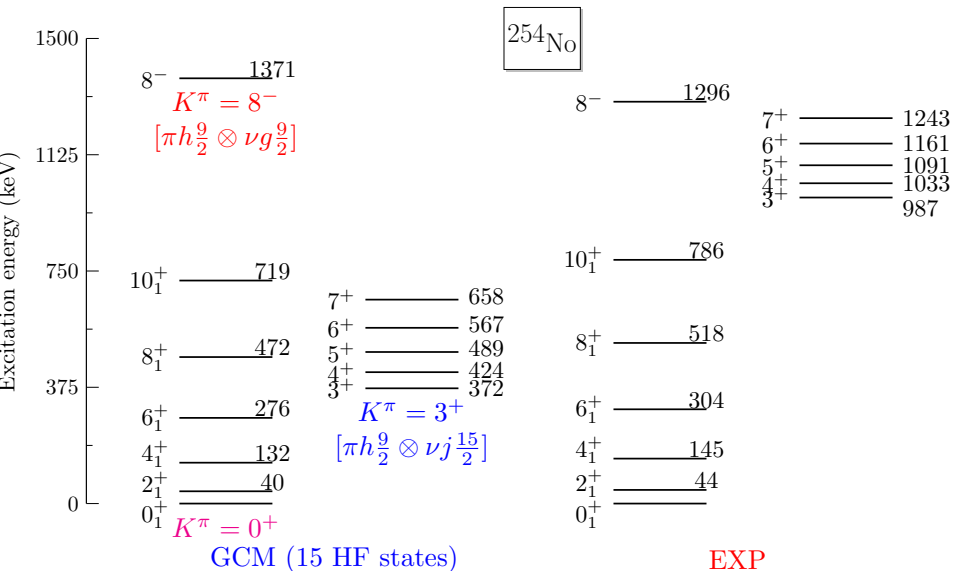
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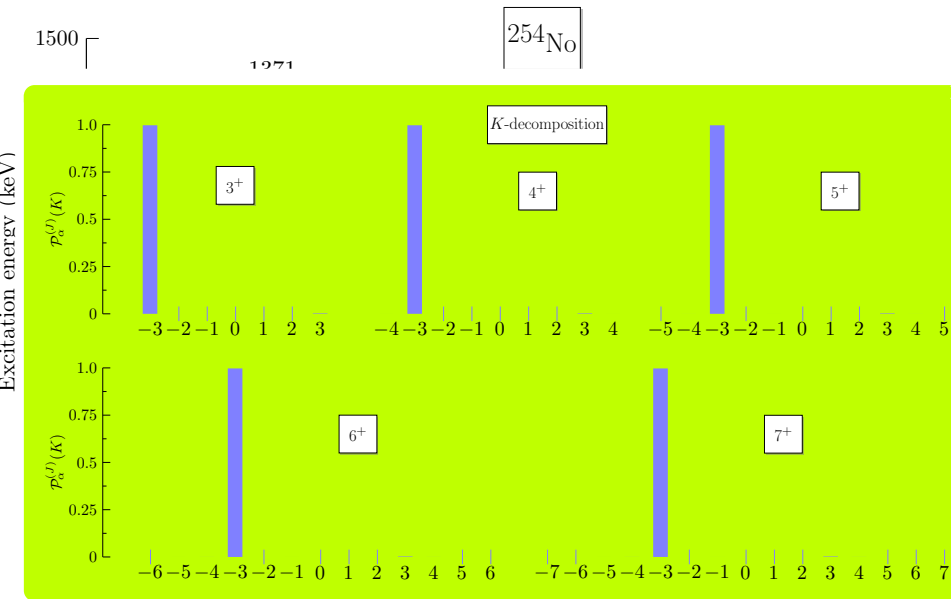


7 <sup>254</sup>102 No<sub>152</sub> : Preliminary calculations with the GCM











7 <sup>254</sup>102 No<sub>152</sub> : Preliminary calculations with the GCM

Occupation numbers:

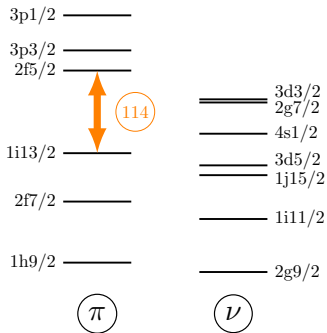
proton orbits	$1h_{9/2}$	$1i_{13/2}$	$2f_{7/2}$	$2f_{5/2}$	$3p_{3/2}$	$3p_{1/2}$	
$0_1^+$	5.66	7.99	3.44	1.58	0.76	0.57	
$8^-$	6.52	7.82	3.28	1.20	0.79	0.39	
$3^+$	6.50	7.98	3.31	1.14	0.72	0.35	
neutron orbits	$1i_{11/2}$	$1j_{15/2}$	$2g_{9/2}$	$2g_{7/2}$	$3d_{5/2}$	$3d_{3/2}$	$4s_{1/2}$
$0_1^+$	7.28	9.67	5.45	1.11	1.16	0.87	0.46
$8^-$	7.29	9.04	6.07	1.12	1.15	0.88	0.45
$3^+$	7.31	9.94	5.43	0.99	1.07	0.83	0.43

### Kuo-Herling interaction:

- <sup>208</sup>Pb<sub>126</sub> core, realistic TBMEs
- $82 \leq Z \leq 126$  shells for proton and  $126 \leq N \leq 184$  for neutrons
- monopole corrections (3N force)

*E. Caurier and F. Nowacki,*

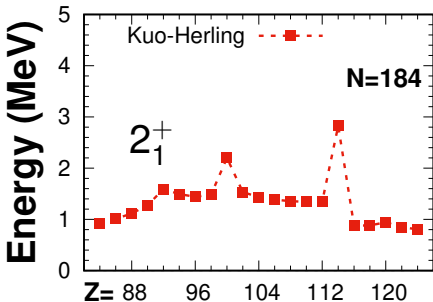
*PRL 87 (2001),072511*



<sup>208</sup>Pb<sub>126</sub>

### Calculations: NATHAN & DFSM

- ◊ Diagonalization within the seniority scheme along the chains of  $N = 126$  and  $N = 184$
- ◊ Deformed Hartree-Fock and Angular momentum projection plus shapes mixing through the GCM: <sup>254</sup>No



**Kuo-Herling interaction:**

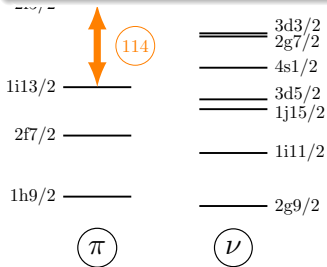
- <sup>208</sup><sub>82</sub>Pb<sub>126</sub> core, realistic TBMEs

**Calculations:** NATHAN & DFSM

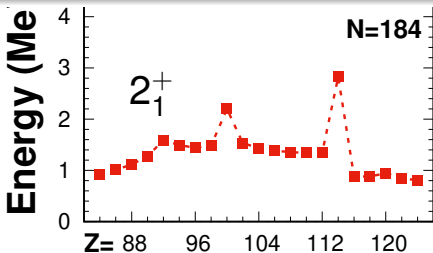
- ◊ Diagonalization within the seniority

**Summary for this part:**

- ◊ First Shell Model interpretation of the superheavy element <sup>254</sup>No
- ◊ Simultaneous description of the rotational band **AND** the isomeric states 8<sup>-</sup> and the 3<sup>+</sup> band
- ◊ Stability predicted at Z = 114



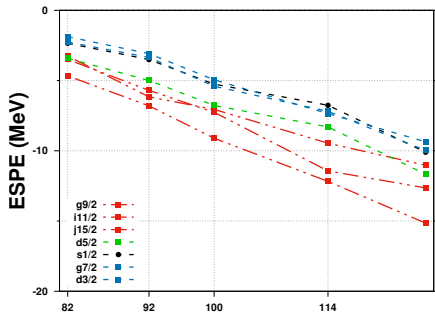
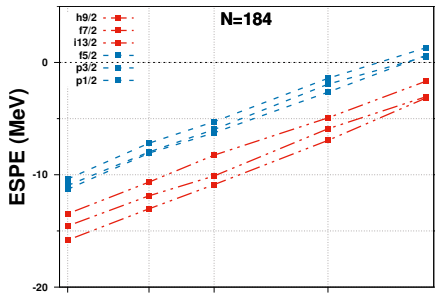
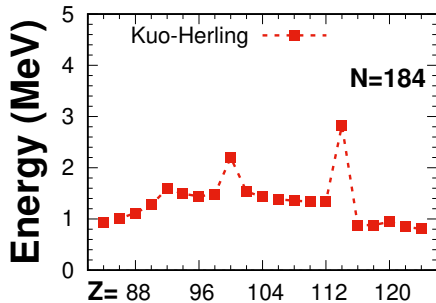
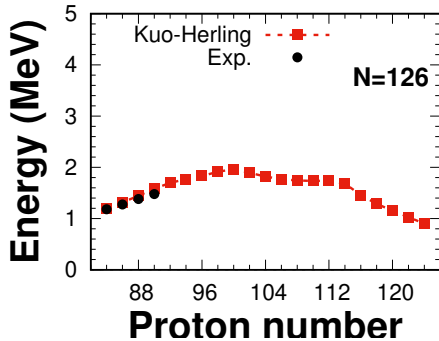
<sup>208</sup><sub>82</sub>Pb<sub>126</sub>



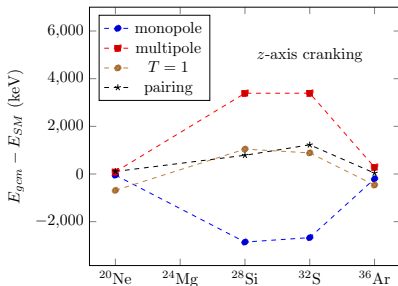
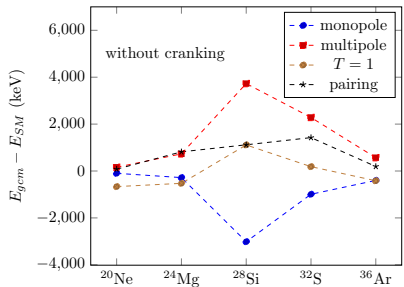
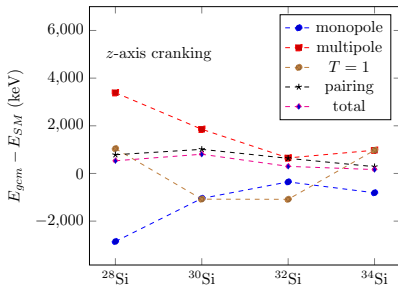
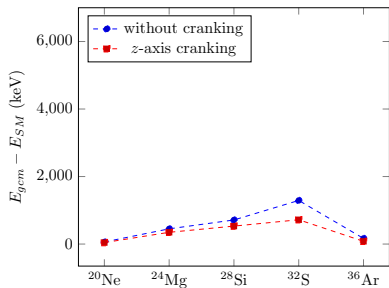
# Summary

- Shell model framework:  
Diagonalization versus Generator Coordinate Method (GCM)
- Systematics calculations of sd shell nuclei
- Collectivity in heavy  $N = Z$  nuclei  
with the effective JUN45.LNPS interaction
- $^{254}\text{No}$ : Preliminary calculations and comparison  
with experimental data (M. Forge talk on Monday)

BACK UP



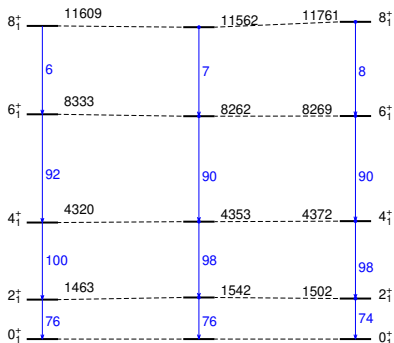
# Systematics in sd nuclei with the USDB interaction



# ② Systematics in sd nuclei with the USDB interaction

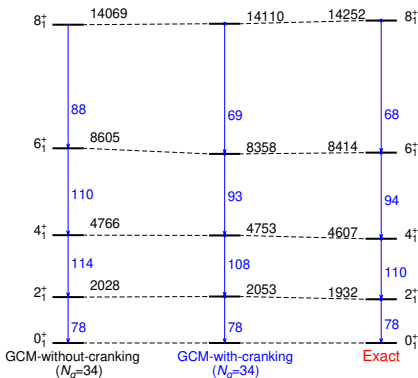
## $^{24}\text{Mg}$

$\Delta E(0_1^+) = +451$  (keV)     $\Delta E_{cr}(0_1^+) = +348$  (keV)



## $^{28}\text{Si}$

$\Delta E(0_1^+) = +714$  (keV)     $\Delta E_{cr}(0_1^+) = +532$  (keV)



### $^{24}\text{Mg}$

$E(2_2^+)$  (keV)

$E(3_1^+)$

$B(E2; 2_2^+ \rightarrow 2_1^+)$  ( $e^2\text{fm}^4$ )

$B(E2; 3_1^+ \rightarrow 2_2^+)$

Exact

4116

5070

14

131

$(\beta, \gamma)$

4059

5140

16

138

$(\beta, \gamma, \langle \hat{J}_z \rangle)$

4121

5199

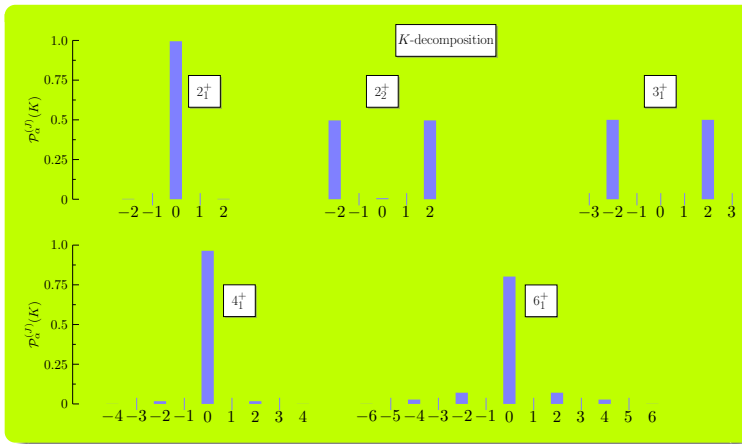
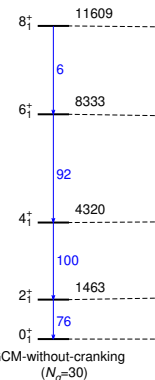
17

137



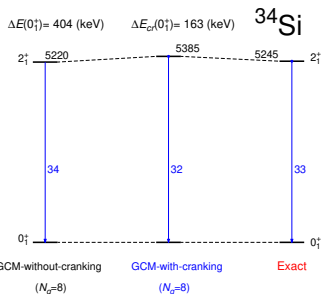
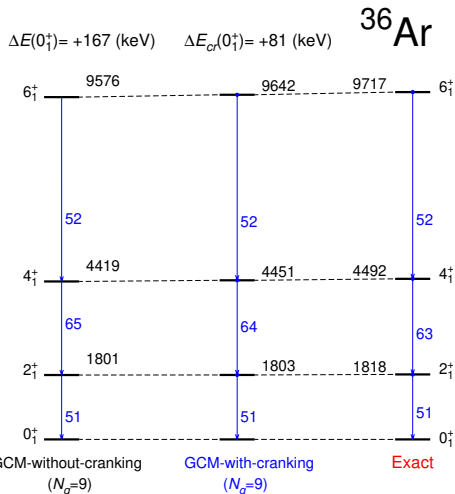
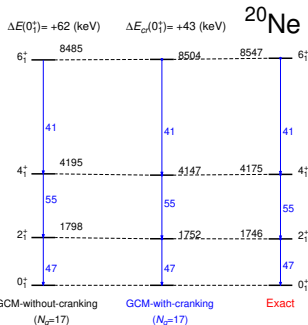
# 2 Systematics in sd nuclei with the USDB interaction

$\Delta E(0_1^+) = +451$  (keV)



$^{24}\text{Mg}$	$E(2_2^+)$ (keV)	$E(3_1^+)$	$B(E2; 2_2^+ \rightarrow 2_1^+)$ ( $e^2\text{fm}^4$ )	$B(E2; 3_1^+ \rightarrow 2_2^+)$
Exact	4116	5070	14	131
$(\beta, \gamma)$	4059	5140	16	138
$(\beta, \gamma, \langle \hat{J}_z \rangle)$	4121	5199	17	137

# 3 Systematics in sd nuclei with the USDB interaction



- Different deformation regimes can be treated with the minimization technique;
- Missing energy of the ground state largest at the middle of the shell  $\sim 350 - 700$  (keV)